

Hamilton-Jacobi Equation for Hamilton's Principal Function

The **Hamilton-Jacobi equation** (HJE) for **Hamilton's Principal Function** is a cornerstone in classical mechanics and provides a bridge to quantum mechanics. Here's an explanation:

1. Context and Background

In classical mechanics, Hamilton's Principal Function, $S(q_1, q_2, \dots, q_n, t)$, arises in the *principle of least action*. It represents the action along a path connecting initial and final configurations in phase space.

- **Action Principle:** The motion of a system can be derived by minimizing the action

$$S = \int L dt,$$

where L is the **Lagrangian** of the system.

- Hamilton's formulation of mechanics involves the Hamiltonian H , which is related to the Lagrangian and describes the energy of the system.

The Hamilton-Jacobi equation is a reformulation of Hamilton's equations, where solving S gives a direct path to solving the equations of motion.

2. Hamilton's Principal Function

S depends on:

- The generalized coordinates q_i ,
- The time t .

It satisfies the condition:

$$\frac{\partial S}{\partial t} = -H\left(q_1, q_2, \dots, q_n; \frac{\partial S}{\partial q_1}, \frac{\partial S}{\partial q_2}, \dots, \frac{\partial S}{\partial q_n}; t\right),$$

where H is the Hamiltonian expressed as a function of generalized coordinates q_i , their conjugate momenta p_i , and time t .

The conjugate momenta are related to S by:

$$p_i = \frac{\partial S}{\partial q_i}.$$

3. The Hamilton-Jacobi Equation

The Hamilton-Jacobi equation can be written as:

$$H\left(q_1, q_2, \dots, q_n; \frac{\partial S}{\partial q_1}, \frac{\partial S}{\partial q_2}, \dots, t\right) + \frac{\partial S}{\partial t} = 0.$$

This is a **partial differential equation** (PDE) for $S(q, t)$. Solving this equation gives the Hamilton's Principal Function, which encapsulates the dynamics of the system.

4. Interpretation

- The **solution** $S(q, t)$ is equivalent to solving the equations of motion for the system.
- Once S is known, the generalized momenta p_i and trajectories $q_i(t)$ can be determined.
- The HJE reduces the problem of solving a set of coupled second-order differential equations (Newton's laws) to solving a first-order PDE.

5. Applications

- **Classical Mechanics:** Direct determination of motion without explicit integration of Hamilton's equations.
- **Quantum Mechanics:** Forms the basis of Schrödinger's equation under certain conditions.
- **Geometrical Optics:** Related to Fermat's principle in optics.
- **General Relativity:** Plays a role in deriving geodesic equations in curved spacetime.

Hamilton's Principal Function and Its Derivation

1. Hamilton's Principal Function as the Action

Hamilton's Principal Function, S , is defined as the **action** along a path in the configuration space:

$$S(q_1, q_2, \dots, q_n, t) = \int_{t_1}^t L(q_i, \dot{q}_i, t) dt,$$

where:

- q_i are the generalized coordinates,
- $\dot{q}_i = \frac{dq_i}{dt}$ are their time derivatives,

- L is the Lagrangian of the system.

This integral depends on the path taken between initial and final points.

2. Total Differential of S

Taking the total differential of S :

$$dS = \frac{\partial S}{\partial t} dt + \sum_i \frac{\partial S}{\partial q_i} dq_i.$$

3. Relation Between S and the Lagrangian

From the definition of S , we have:

$$S = \int_{t_1}^t L dt.$$

Differentiating S with respect to time:

$$\frac{\partial S}{\partial t} = L.$$

4. Relation Between S and Conjugate Momentum

The conjugate momentum p_i is defined as:

$$p_i = \frac{\partial L}{\partial \dot{q}_i}.$$

Using the principle of least action, the variation of the action S leads to Hamilton's equations. To find the connection between p_i and S , observe the dependence of S on the coordinates q_i . From the total differential of S :

$$dS = \sum_i \frac{\partial S}{\partial q_i} dq_i + \frac{\partial S}{\partial t} dt.$$

Comparing this with the canonical Hamiltonian formulation, we identify:

$$p_i = \frac{\partial S}{\partial q_i}.$$

5. Hamilton-Jacobi Equation

From Hamilton's equations, the Hamiltonian H is related to S as:

$$H(q_i, p_i, t) = \sum_i p_i \dot{q}_i - L.$$

Since $p_i = \frac{\partial S}{\partial q_i}$, substituting into the expression for H and using $\frac{\partial S}{\partial t} = -H$, we get:

$$\frac{\partial S}{\partial t} + H\left(q_i, \frac{\partial S}{\partial q_i}, t\right) = 0.$$

This is the Hamilton-Jacobi equation, with S encapsulating the dynamics.

Summary of Key Relationships

1. Hamilton's Principal Function:

$$S = \int L dt.$$

2. Conjugate momentum:

$$p_i = \frac{\partial S}{\partial q_i}.$$

3. Hamilton-Jacobi equation:

$$\frac{\partial S}{\partial t} + H\left(q_i, \frac{\partial S}{\partial q_i}, t\right) = 0.$$

Why is Hamilton's Principal Function Independent of \dot{q}_i When L Depends on \dot{q}_i ?

Hamilton's Principal Function $S(q_i, t)$ is independent of \dot{q}_i (the generalized velocities), even though the Lagrangian $L(q_i, \dot{q}_i, t)$ depends on \dot{q}_i . Here's the reasoning:

1. Definition of S

Hamilton's Principal Function is defined as:

$$S(q_i, t) = \int_{t_1}^t L(q_i, \dot{q}_i, t) dt.$$

This integral is evaluated along a specific trajectory in the configuration space that satisfies the equations of motion (derived from the principle of least action). Thus, S is not a general function of q_i , \dot{q}_i , and t , but instead depends only on the final generalized coordinates q_i and time t , after the trajectory is determined.

2. Trajectory Dependence on \dot{q}_i

The Lagrangian $L(q_i, \dot{q}_i, t)$ explicitly depends on \dot{q}_i . However:

- The trajectory of the system is uniquely determined by the equations of motion (Euler-Lagrange equations) and the boundary conditions.
- Along this trajectory, \dot{q}_i is no longer an independent variable; it is a function of q_i , t , and possibly initial conditions.

Thus, when integrating L along the trajectory, the dependence on \dot{q}_i is “absorbed” into the dependence of S on q_i and t .

3. Reduction to Canonical Coordinates

Hamilton's Principal Function S is constructed in such a way that it serves as a generator of the canonical transformation connecting the generalized coordinates q_i and their conjugate momenta p_i . The conjugate momenta are defined as:

$$p_i = \frac{\partial L}{\partial \dot{q}_i}.$$

Since \dot{q}_i appears only implicitly in S via $p_i = \frac{\partial S}{\partial q_i}$, S itself does not depend explicitly on \dot{q}_i .

4. Role of \dot{q}_i in the Variational Principle

The dependence of L on \dot{q}_i ensures that the equations of motion (Euler-Lagrange equations) can be derived from the action \mathcal{S} :

$$\mathcal{S} = \int_{t_1}^{t_2} L(q_i, \dot{q}_i, t) dt.$$

Once the equations of motion are solved, the explicit dependence on \dot{q}_i vanishes because the velocities are determined by the trajectory $q_i(t)$.

5. Summary

- **Lagrangian Dependence:** L depends on \dot{q}_i because it describes the kinetic and potential energies of the system in terms of both positions q_i and velocities \dot{q}_i .
- **Hamilton's Principal Function Dependence:** $S(q_i, t)$ is the integrated action along a physical trajectory. As such, it depends only on the generalized coordinates q_i and time t , not on the generalized velocities \dot{q}_i , which are implicitly accounted for in the trajectory.

Derivation of the Hamilton-Jacobi Equation

We derive the Hamilton-Jacobi Equation (HJE) step by step, explicitly replacing L with $\frac{dS}{dt}$.

1. Hamilton's Principal Function

Hamilton's Principal Function $S(q_i, t)$ is defined as:

$$S(q_i, t) = \int L dt,$$

where $L(q_i, \dot{q}_i, t)$ is the Lagrangian. The total time derivative of S is:

$$\frac{dS}{dt} = L(q_i, \dot{q}_i, t).$$

2. Hamiltonian Definition

The Hamiltonian is defined as:

$$H = \sum_i p_i \dot{q}_i - L,$$

where:

- $p_i = \frac{\partial L}{\partial \dot{q}_i}$ is the conjugate momentum,
- $\dot{q}_i = \frac{dq_i}{dt}$ is the generalized velocity.

Now, replace L with $\frac{dS}{dt}$ in this definition. The total time derivative of S is:

$$\frac{dS}{dt} = \frac{\partial S}{\partial t} + \sum_i \frac{\partial S}{\partial q_i} \dot{q}_i.$$

Substituting $\frac{dS}{dt}$ into the Hamiltonian, we get:

$$H = \sum_i p_i \dot{q}_i - \left(\frac{\partial S}{\partial t} + \sum_i \frac{\partial S}{\partial q_i} \dot{q}_i \right).$$

3. Simplification Using $p_i = \frac{\partial S}{\partial q_i}$

Since $p_i = \frac{\partial S}{\partial q_i}$, the term $\sum_i p_i \dot{q}_i$ cancels with $\sum_i \frac{\partial S}{\partial q_i} \dot{q}_i$. This leaves:

$$H = -\frac{\partial S}{\partial t}.$$

4. Substituting into the Hamiltonian

The Hamiltonian H is a function of q_i , p_i , and t . Substituting $p_i = \frac{\partial S}{\partial q_i}$, we express H as:

$$H \left(q_i, \frac{\partial S}{\partial q_i}, t \right).$$

Using the relation $H = -\frac{\partial S}{\partial t}$, we obtain:

$$\frac{\partial S}{\partial t} + H \left(q_i, \frac{\partial S}{\partial q_i}, t \right) = 0.$$

5. Final Form of the Hamilton-Jacobi Equation

The final form of the Hamilton-Jacobi Equation is:

$$\frac{\partial S}{\partial t} + H \left(q_i, \frac{\partial S}{\partial q_i}, t \right) = 0,$$

where:

- $S(q_i, t)$ is Hamilton's Principal Function,
- $p_i = \frac{\partial S}{\partial q_i}$.

6. Summary of Key Steps

1. Start with $S(q_i, t) = \int L dt$.
2. Replace $L = \frac{dS}{dt} = \frac{\partial S}{\partial t} + \sum_i \frac{\partial S}{\partial q_i} \dot{q}_i$.
3. Substitute into the Hamiltonian definition:

$$H = \sum_i p_i \dot{q}_i - L.$$

4. Cancel terms using $p_i = \frac{\partial S}{\partial q_i}$.
5. Arrive at the Hamilton-Jacobi Equation:

$$\frac{\partial S}{\partial t} + H\left(q_i, \frac{\partial S}{\partial q_i}, t\right) = 0.$$

Harmonic Oscillator Problem Using the Hamilton-Jacobi Method

1. The Harmonic Oscillator Hamiltonian

The Hamiltonian of a one-dimensional harmonic oscillator is given by:

$$H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 q^2,$$

where:

- q : Generalized coordinate,
- p : Conjugate momentum,
- m : Mass of the oscillator,
- ω : Angular frequency of the oscillator.

2. Hamilton-Jacobi Equation

The Hamilton-Jacobi equation (HJE) for Hamilton's Principal Function $S(q, t)$ is:

$$\frac{\partial S}{\partial t} + H\left(q, \frac{\partial S}{\partial q}\right) = 0.$$

Substitute H into the equation:

$$\frac{\partial S}{\partial t} + \frac{1}{2m} \left(\frac{\partial S}{\partial q}\right)^2 + \frac{1}{2}m\omega^2 q^2 = 0.$$

3. Separation of Variables

Assume a solution of the form:

$$S(q, t) = W(q) - Et,$$

where $W(q)$ is the time-independent part of S , and E is the energy of the system.

Substituting into the HJE:

$$-\frac{\partial S}{\partial t} + \frac{1}{2m} \left(\frac{\partial W}{\partial q} \right)^2 + \frac{1}{2} m \omega^2 q^2 = 0.$$

Simplify:

$$E = \frac{1}{2m} \left(\frac{\partial W}{\partial q} \right)^2 + \frac{1}{2} m \omega^2 q^2.$$

Rearranging gives:

$$\left(\frac{\partial W}{\partial q} \right)^2 = 2mE - m^2 \omega^2 q^2.$$

4. Solve for $W(q)$

Take the square root of both sides:

$$\frac{\partial W}{\partial q} = \pm \sqrt{2mE - m^2 \omega^2 q^2}.$$

Integrate:

$$W(q) = \int \pm \sqrt{2mE - m^2 \omega^2 q^2} dq.$$

Perform the integration (using standard techniques for a quadratic under the square root):

$$W(q) = \pm \frac{m\omega}{2} \left(q \sqrt{\frac{2E}{m\omega^2} - q^2} + \frac{1}{\omega} \arcsin \left(\sqrt{\frac{m\omega^2}{2E}} q \right) \right).$$

5. Hamilton's Characteristic Function and Motion

The total Hamilton's Principal Function is:

$$S(q, t) = W(q) - Et.$$

From $W(q)$, the momentum is:

$$p = \frac{\partial S}{\partial q} = \frac{\partial W}{\partial q} = \pm \sqrt{2mE - m^2 \omega^2 q^2}.$$

6. Action-Angle Variables

The Hamilton-Jacobi method naturally introduces **action-angle variables** for periodic systems like the harmonic oscillator. The action J is given by:

$$J = \oint p dq.$$

Substitute $p = \sqrt{2mE - m^2\omega^2q^2}$:

$$J = \int_{-q_{\max}}^{q_{\max}} \sqrt{2mE - m^2\omega^2q^2} dq.$$

Perform the integral (this is the area of an ellipse in phase space):

$$J = \frac{E}{\omega}.$$

The angular frequency ω is then directly related to the action.

Summary of Results

1. The energy of the harmonic oscillator:

$$E = \frac{1}{2}m\omega^2q^2 + \frac{p^2}{2m}.$$

2. The Hamilton-Jacobi equation is solved by:

$$S(q, t) = W(q) - Et,$$

where $W(q)$ is derived as above.

3. The system can be described in terms of action-angle variables for periodic motion.

Logic Behind the Variable Separable Form in the Hamilton-Jacobi Method

1. The Logic Behind the Separable Form

The chosen separable form for $S(q, t)$:

$$S(q, t) = W(q) - Et,$$

is guided by the following considerations:

a. The Structure of the Hamilton-Jacobi Equation

The Hamilton-Jacobi equation is:

$$\frac{\partial S}{\partial t} + H\left(q, \frac{\partial S}{\partial q}\right) = 0.$$

- H , the Hamiltonian, often depends only on the generalized coordinate q , the conjugate momentum p , and possibly t .
- A natural way to handle the explicit dependence on t is to separate it from the spatial variables.

b. Energy Conservation

For systems with a conserved energy E (time-independent Hamiltonians), $S(q, t)$ can often be expressed as:

$$S(q, t) = W(q) - Et,$$

where:

- $W(q)$ is the time-independent part of S ,
- E represents the total energy, acting as a separation constant.

This separation reflects the symmetry of the system: the time evolution is independent of the spatial configuration q .

2. Can Other Variable Separable Forms Be Used?

Yes, other separable forms can be used, but the specific form depends on:

- The structure of the Hamiltonian,
- The nature of the system's constraints and symmetries.

a. Time-Independent Systems

If H is time-independent, the form $S(q, t) = W(q) - Et$ is natural because t separates cleanly as the Hamiltonian itself defines E . Any other form would likely complicate the PDE unnecessarily.

b. Time-Dependent Systems

For time-dependent Hamiltonians $H(q, p, t)$, separable forms such as:

$$S(q, t) = W(q) + F(t),$$

might be more appropriate, where $F(t)$ is determined based on the time-dependence of H . This is seen, for example, in systems with external time-dependent forces.

c. Multidimensional Systems

In multidimensional systems with separable coordinates q_1, q_2, \dots , S can be separated into components:

$$S(q_1, q_2, \dots, t) = \sum_i W_i(q_i) - Et,$$

where $W_i(q_i)$ corresponds to the motion in the i -th coordinate.

d. Arbitrary Forms

For systems with more complex constraints, other separable forms might emerge, but they would generally follow from:

- The symmetries of the system (e.g., spherical, cylindrical),
- The conserved quantities that allow separation (e.g., energy, angular momentum).

3. Why the Standard Form Works for the Harmonic Oscillator

In the harmonic oscillator:

$$H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 q^2,$$

- H is time-independent,
- Energy E is conserved.

The form $S(q, t) = W(q) - Et$ leverages these properties:

- t is treated independently due to energy conservation,
- q and p are connected through $W(q)$ and the HJE.

4. Why Not Use Arbitrary Separable Forms?

Other separable forms may be theoretically possible, but:

- They might introduce unnecessary complexity in solving the HJE.
- The standard form $S(q, t) = W(q) - Et$ is chosen because it aligns directly with conserved quantities (like E) and simplifies the PDE into an ODE for $W(q)$.

Summary

- The standard form $S(q, t) = W(q) - Et$ is motivated by the time-independence of the Hamiltonian and energy conservation.
- Other separable forms are possible but are typically tailored to the system's symmetries, constraints, or time dependence.
- The choice of form is guided by the goal of simplifying the Hamilton-Jacobi equation while reflecting the physical properties of the system.

Hamilton's Characteristic Function and Principal Function

Hamilton's Characteristic Function and Principal Function are closely related but have distinct purposes. Here's a detailed explanation:

1. Hamilton's Principal Function ($S(q_i, t)$)

Hamilton's Principal Function $S(q_i, t)$ is defined as:

$$S(q_i, t) = \int L dt,$$

where L is the **Lagrangian** of the system.

Key Characteristics:

- $S(q_i, t)$ depends explicitly on the generalized coordinates q_i and the time t .
- It solves the **time-dependent Hamilton-Jacobi Equation (HJE)**:

$$\frac{\partial S}{\partial t} + H\left(q_i, \frac{\partial S}{\partial q_i}, t\right) = 0,$$

where H is the Hamiltonian.

- S describes the full dynamics of the system, including time evolution.

2. Hamilton's Characteristic Function ($W(q_i)$)

Hamilton's Characteristic Function $W(q_i)$ is the time-independent version of $S(q_i, t)$. It is used for systems where the Hamiltonian is **time-independent**.

Definition:

Hamilton's Characteristic Function $W(q_i)$ solves the **time-independent Hamilton-Jacobi Equation**:

$$H\left(q_i, \frac{\partial W}{\partial q_i}\right) = E,$$

where:

- E is the total energy of the system,
- $H(q_i, p_i)$ is the Hamiltonian.

Key Characteristics:

- $W(q_i)$ depends only on the generalized coordinates q_i and constants of motion (like E).
- It does not involve time t explicitly.
- W describes the geometry of trajectories in phase space rather than their evolution in time.

3. Relation Between S and W

Hamilton's Principal Function $S(q_i, t)$ and Hamilton's Characteristic Function $W(q_i)$ are related for systems with time-independent Hamiltonians. In such cases:

$$S(q_i, t) = W(q_i) - Et,$$

where:

- $W(q_i)$ is Hamilton's Characteristic Function,
- E is the total energy,
- t is time.

This shows that S is a combination of $W(q_i)$ and the time-dependent term $-Et$.

4. When to Use Each Function

Function	When to Use
Hamilton's Principal Function (S)	For solving problems with time-dependent Hamiltonians ($H = H(q_i, p_i, t)$).
Hamilton's Characteristic Function (W)	For time-independent Hamiltonians ($H = H(q_i, p_i)$).

5. Example: Harmonic Oscillator

Hamiltonian:

$$H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 q^2.$$

Time-Independent HJE (for W):

$$\frac{1}{2m} \left(\frac{\partial W}{\partial q} \right)^2 + \frac{1}{2}m\omega^2 q^2 = E.$$

Solving this gives $W(q)$, the characteristic function:

$$W(q) = \pm \int \sqrt{2mE - m^2\omega^2 q^2} dq.$$

Time-Dependent HJE (for S):

Using $S(q, t) = W(q) - Et$, the principal function is:

$$S(q, t) = \pm \int \sqrt{2mE - m^2\omega^2 q^2} dq - Et.$$

6. Summary

- **Hamilton's Principal Function (S):** Time-dependent, describes the full dynamics of the system.
- **Hamilton's Characteristic Function (W):** Time-independent, focuses on the geometry of trajectories.

They are related via:

$$S(q_i, t) = W(q_i) - Et,$$

for time-independent systems.

Separation of Variables in the Hamilton-Jacobi Equation

The Hamilton-Jacobi Equation (HJE) is:

$$\frac{\partial S}{\partial t} + H \left(q_i, \frac{\partial S}{\partial q_i}, t \right) = 0,$$

where:

- $S(q_i, t)$ is Hamilton's Principal Function,

- H is the Hamiltonian.

For **time-independent systems**, the HJE simplifies to:

$$H\left(q_i, \frac{\partial S}{\partial q_i}\right) = E,$$

where E is the total energy of the system.

1. Separation of Variables in the HJE

Separation of variables is a method to solve the HJE by assuming S can be expressed as a sum of functions, each depending on a single variable.

Assumption:

For a system with n generalized coordinates q_1, q_2, \dots, q_n , assume:

$$S(q_1, q_2, \dots, q_n, t) = W(q_1, q_2, \dots, q_n) - Et,$$

where:

- $W(q_1, q_2, \dots, q_n)$ is Hamilton's Characteristic Function,
- E is the total energy.

Further assume W can be written as a sum:

$$W(q_1, q_2, \dots, q_n) = \sum_{i=1}^n W_i(q_i),$$

where W_i depends only on q_i .

Substitution into the Time-Independent HJE:

Substitute $S = W - Et$ into the HJE:

$$H\left(q_i, \frac{\partial S}{\partial q_i}\right) = E.$$

If H is separable, this becomes:

$$\sum_{i=1}^n H_i\left(q_i, \frac{\partial W_i}{\partial q_i}\right) = E,$$

where H_i depends only on q_i and $\frac{\partial W_i}{\partial q_i}$.

Separation:

Each term H_i is equated to a constant α_i , such that:

$$H_i \left(q_i, \frac{\partial W_i}{\partial q_i} \right) = \alpha_i,$$

and:

$$\sum_{i=1}^n \alpha_i = E.$$

This transforms the HJE into n simpler, independent equations.

2. Cyclic (Ignorable) Coordinates

A **cyclic coordinate** (or ignorable coordinate) is a coordinate q_j that does not appear explicitly in the Hamiltonian H . This property simplifies the HJE.

Case of a Cyclic Coordinate:

If q_j is cyclic:

$$H = H(q_1, q_2, \dots, p_j, \dots),$$

and p_j (the conjugate momentum) is constant:

$$p_j = \frac{\partial S}{\partial q_j} = \text{constant}.$$

Impact on Separability:

For a cyclic coordinate q_j , the corresponding term in $W(q_i)$ is linear in q_j :

$$W_j(q_j) = p_j q_j.$$

This simplifies the separation of variables. The total function W becomes:

$$W(q_1, q_2, \dots) = p_j q_j + \sum_{i \neq j} W_i(q_i).$$

3. Example: Particle in a Central Potential**Hamiltonian:**

$$H = \frac{p_r^2}{2m} + \frac{p_\phi^2}{2mr^2} + V(r).$$

Cyclic Coordinate:

- ϕ is cyclic because it does not appear in $V(r)$.
- Thus:

$$p_\phi = \frac{\partial S}{\partial \phi} = \text{constant}.$$

Separation of Variables:

Assume:

$$S(r, \phi, t) = W_r(r) + W_\phi(\phi) - Et.$$

Substitute into the HJE and separate $W_r(r)$ and $W_\phi(\phi)$:

$$\frac{p_r^2}{2m} + \frac{p_\phi^2}{2mr^2} + V(r) = E.$$

This leads to:

1. An equation for $W_\phi(\phi) = p_\phi \phi$,
2. A radial equation for $W_r(r)$.

4. Summary of Key Steps

1. Separability in the HJE:

- Assume $S(q_i, t) = W(q_i) - Et$.
- Further assume $W(q_i) = \sum_{i=1}^n W_i(q_i)$.
- Substitute into the HJE and separate into n independent equations.

2. Cyclic Coordinates:

- A cyclic coordinate q_j does not appear in H , making p_j constant.
- For a cyclic coordinate, $W_j(q_j) = p_j q_j$, simplifying the solution.

3. Applications:

- This method is widely used for systems with symmetry, such as central force problems, harmonic oscillators, and planetary motion.

Action-Angle Variables

The **action-angle variables** are a set of canonical coordinates used to describe systems with periodic or quasi-periodic motion, such as harmonic oscillators, planetary orbits, or particles in central potentials. These variables simplify the dynamics of such systems.

1. Key Concepts

(a) Action Variable (J_i)

- The **action variable** is a conserved quantity associated with periodic motion.

- It is defined as the integral of the conjugate momentum p_i over one complete cycle of its periodic motion:

$$J_i = \oint p_i dq_i,$$

where q_i is the generalized coordinate and p_i is the conjugate momentum.

- J_i is constant for integrable systems.

(b) Angle Variable (θ_i)

- The **angle variable** represents the phase of the motion and changes linearly with time.
- It evolves as:

$$\theta_i(t) = \theta_{i0} + \omega_i t,$$

where $\omega_i = \frac{\partial H}{\partial J_i}$ is the angular frequency, and H is the Hamiltonian.

2. Why Use Action-Angle Variables?

- **Simplification:** For integrable systems, the equations of motion in action-angle variables are simple. The Hamiltonian H depends only on the action variables:

$$H = H(J_1, J_2, \dots).$$

The angle variables evolve linearly:

$$\dot{\theta}_i = \omega_i = \frac{\partial H}{\partial J_i}.$$

- **Periodic Systems:** These variables are ideal for systems where motion is periodic or quasi-periodic, as the angle variable θ_i captures the periodicity.
- **Canonical Transformation:** The transformation from (q_i, p_i) to (J_i, θ_i) is canonical, preserving the structure of Hamilton's equations.

3. How Are They Defined?

(a) Action Variable (J_i)

The action variable J_i is the area enclosed by the trajectory in phase space for the i -th degree of freedom:

$$J_i = \oint p_i dq_i.$$

(b) Angle Variable (θ_i)

The angle variable θ_i parameterizes the position within the periodic trajectory. It is defined such that:

$$\theta_i = \frac{\partial W}{\partial J_i},$$

where W is the generating function of the canonical transformation to action-angle coordinates.

4. Example: Harmonic Oscillator

Hamiltonian:

$$H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 q^2.$$

(a) Action Variable (J):

The trajectory in phase space is an ellipse. The action variable is the area enclosed by this ellipse:

$$J = \oint p dq.$$

Using energy conservation:

$$p = \sqrt{2mE - m^2\omega^2 q^2},$$

we compute J :

$$J = \int_{-q_{\max}}^{q_{\max}} \sqrt{2mE - m^2\omega^2 q^2} dq = \frac{E}{\omega}.$$

(b) Angle Variable (θ):

The angle variable evolves linearly with time:

$$\theta(t) = \omega t + \theta_0,$$

where ω is the angular frequency.

5. Applications

- **Perturbation Theory:** Used to study small deviations from integrable systems.
- **Celestial Mechanics:** Describes planetary orbits in a central force field.
- **Quantum Mechanics:** Quantization of the action variable J_i leads to quantum conditions:

$$J_i = n_i h, \quad n_i \in \mathbf{Z}.$$

6. Summary

- **Action Variables (J_i):** Conserved quantities that are integrals of motion for periodic systems.
- **Angle Variables (θ_i):** Periodic variables that describe the phase of the motion.
- Action-angle variables simplify the study of periodic and quasi-periodic systems by reducing the dynamics to simple linear evolution in the angle variables.